

Particle Interactions



Heavy Particle Collisions



- Charged particles interact in matter.
 - Ionization and excitation of atoms
 - Nuclear interactions rare
- Electrons can lose most of their energy in a single collision with an atomic electron.
- Heavier charged particles lose a small fraction of their energy to atomic electrons with each collision.

Energy Transfer

- Assume an elastic collision.
 - One dimension
 - Moving particle M , V
 - Initial energy $E = \frac{1}{2} MV^2$
 - Electron mass m
 - Outgoing velocities V_f , v_f
- This gives a maximum energy transfer Q_{max} .

$$MV = MV_f + mv_f$$

$$\frac{1}{2} MV^2 = \frac{1}{2} MV_f^2 + \frac{1}{2} mv_f^2$$



$$V_f = \frac{M - m}{M + m} V$$

$$\begin{aligned} Q_{max} &= \frac{1}{2} MV^2 - \frac{1}{2} MV_f^2 \\ &= \frac{4mM}{(M + m)^2} \left(\frac{1}{2} MV^2 \right) \cong \frac{4mE}{M} \end{aligned}$$

Relativistic Energy Transfer

- At high energy relativistic effects must be included.

$$Q_{\max} = \frac{2\gamma^2 m V^2}{1 + 2\gamma m / M + m^2 / M^2}$$

- This reduces for heavy particles at low γ , $\gamma m / M \ll 1$.

$$Q_{\max} \cong 2\gamma^2 m V^2 = 2\gamma^2 \beta^2 m c^2$$

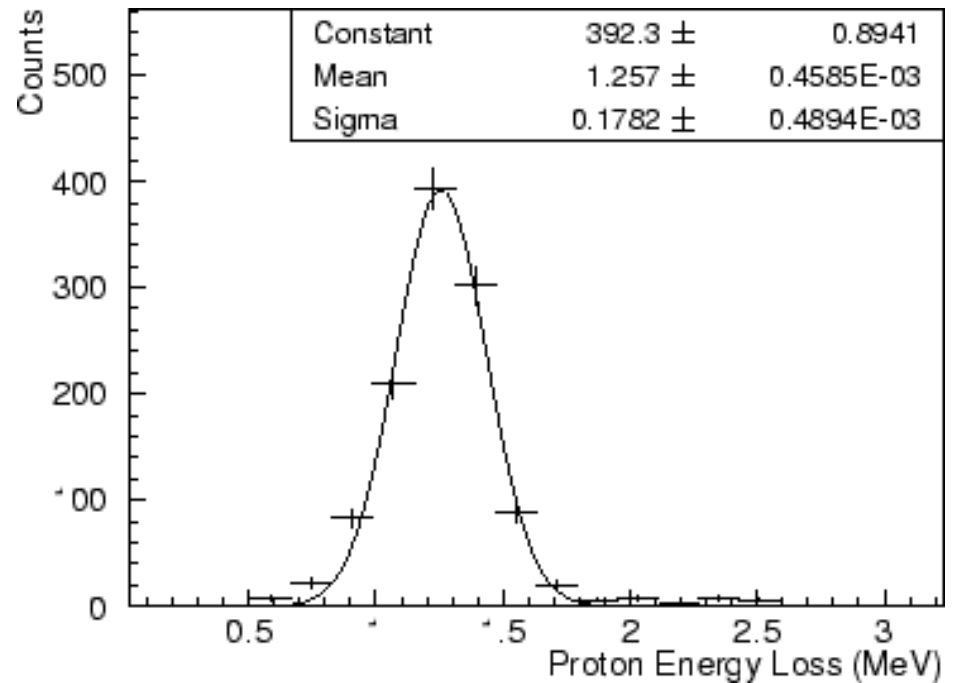
$$Q_{\max} \cong 2(\gamma^2 - 1) m c^2$$

Typical Problem

- Calculate the maximum energy a 100-MeV pion can transfer to an electron.
- $m_{\pi} = 139.6 \text{ MeV} = 280 m_e$
 - problem is relativistic
- $\gamma = (K + m_{\pi}) / m_{\pi} = 2.4$
- $Q_{\max} = 5.88 \text{ MeV}$

Protons in Silicon

- The MIDAS detector measured proton energy loss.
 - 125 MeV protons
 - Thin wafer of silicon



Linear Energy Transfer

- Charged particles experience multiple interactions passing through matter.
 - Integral of individual collisions
 - Probability $P(Q)$ of energy transfer Q

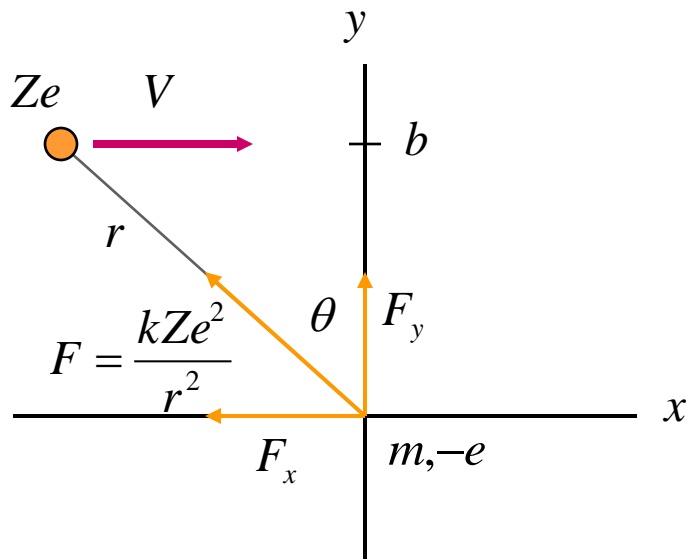
$$\bar{Q} = \int_{Q_{\min}}^{Q_{\max}} QP(Q)dQ$$

- Rate of loss is the *stopping power* or *LET* or dE/dx .
 - Use probability of a collision μ (cm^{-1})

$$-\frac{dE}{dx} = \mu\bar{Q} = \mu \int_{Q_{\min}}^{Q_{\max}} QP(Q)dQ$$

Energy Loss

- The stopping power can be derived semiclassically.
 - Heavy particle Ze, V
 - Impact parameter b



- Calculate the impulse and energy to the electron.

$$p = \int F_y dt = \int \frac{kZe^2}{r^2} \frac{b}{r} dt$$

$$p = kZe^2 \int \frac{b}{(b^2 + V^2 t^2)^{3/2}} dt$$

$$p = \frac{2kZe^2}{Vb}$$

$$Q = \frac{p^2}{2m} = \frac{2k^2 Z^2 e^4}{mV^2 b^2}$$

Impact Parameter

- Assume a uniform density of electrons.
 - n per unit volume
 - Thickness dx
- Consider an impact parameter range b to $b+db$
 - Integrate over range of b
 - Equivalent to range of Q

- Find the number of electrons.

$$2\pi n b(db) dx$$

- Now find the energy loss per unit distance.

$$-\frac{dE}{dx} = 2\pi n \int Q b db = \frac{4\pi n k^2 Z^2 e^4}{mV^2} \int \frac{db}{b}$$

$$-\frac{dE}{dx} = \frac{4\pi n k^2 Z^2 e^4}{mV^2} \ln \frac{b_{\max}}{b_{\min}}$$

Stopping Power

- The impact parameter is related to characteristic frequencies.
- Compare the maximum b to the orbital frequency f .
 - $b/V < 1/f$
 - $b_{\max} = V/f$
- Compare the minimum b to the de Broglie wavelength.
 - $b_{\min} = h / mV$

- The classical Bohr stopping power is

$$-\frac{dE}{dx} = \frac{4\pi n k^2 Z^2 e^4}{mV^2} \ln \frac{mV^2}{hf}$$

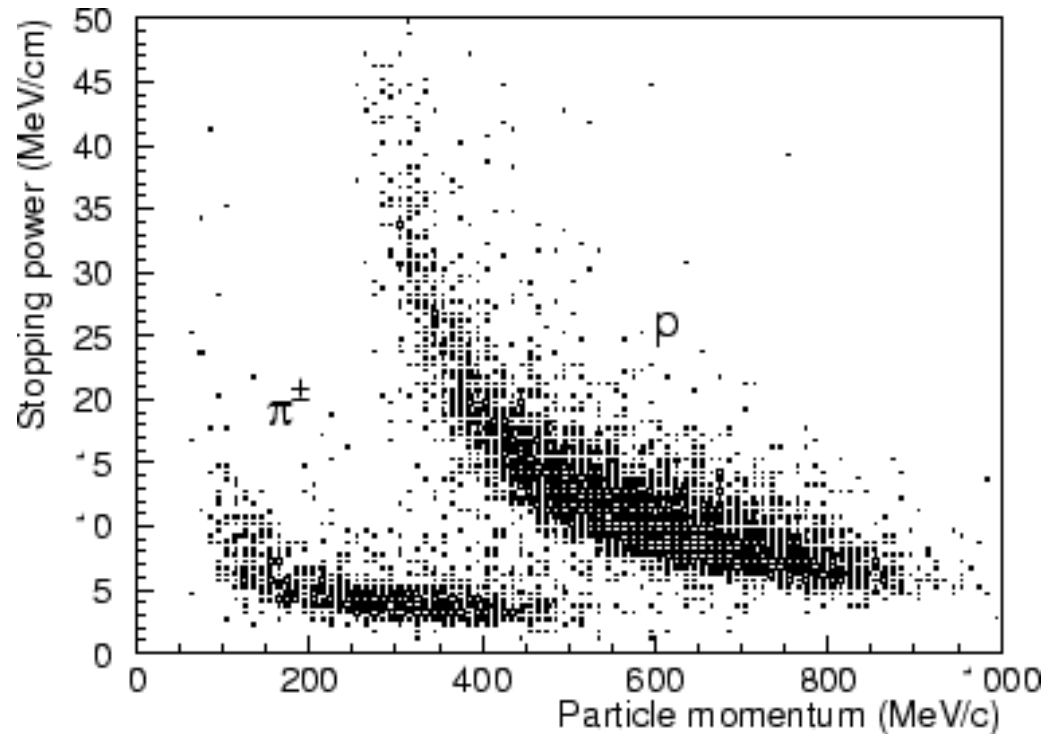
Bethe Formula

- A complete treatment of stopping power requires relativistic quantum mechanics.
 - Include speed b
 - Material dependent excitation energy

$$-\frac{dE}{dx} = \frac{4\pi n k^2 Z^2 e^4}{mc^2 \beta^2} \left[\ln \frac{2mc^2 \beta^2}{I(1-\beta^2)} - \beta^2 \right]$$

Silicon Stopping Power

- Protons and pions behave differently in matter
 - Different mass
 - Energy dependent



Range

- Range is the distance a particle travels before coming to rest.
- Stopping power represents energy per distance.
 - Range based on energy
- Use Bethe formula with term that only depends on speed
 - Numerically integrated
 - Used for mass comparisons

$$R(K) = \int_0^K \left(-\frac{dE}{dx} \right)^{-1} dE$$

$$R(K) = \frac{1}{Z^2} \int_0^K \frac{dE}{G(\beta)}$$

$$R(\beta) = \frac{1}{Z^2} \int_0^\beta \frac{Mg(\beta)d\beta}{G(\beta)} = \frac{M}{Z^2} f(\beta)$$

Alpha Penetration

Typical Problem

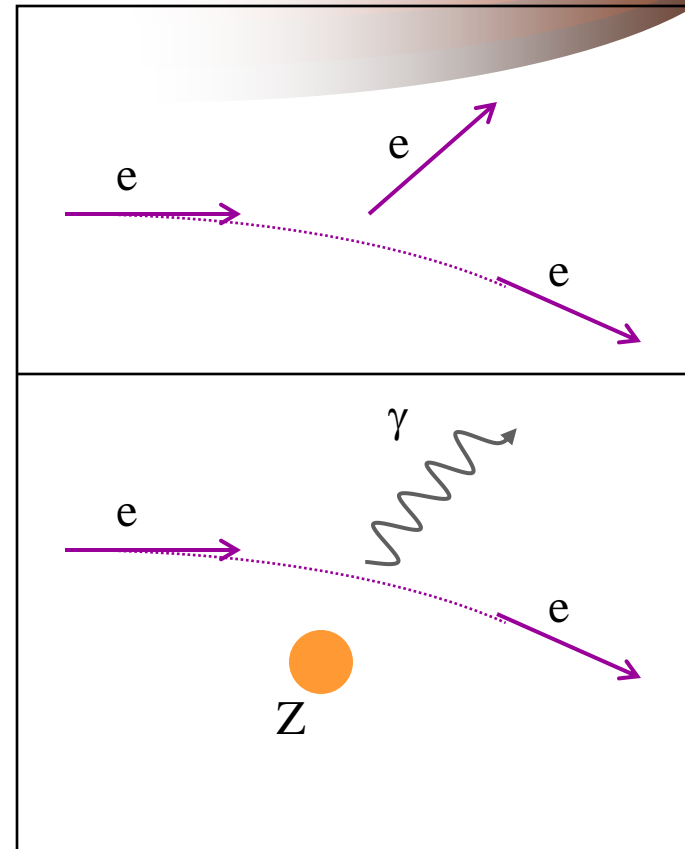
- Part of the radon decay chain includes a 7.69 MeV alpha particle. What is the range of this particle in soft tissue?
- Use proton mass and charge equal to 1.
 - $M_\alpha = 4, Z^2 = 4$

$$R_\alpha(\beta) = \frac{M}{Z^2} R_p(\beta) = R_p(\beta)$$

- Equivalent energy for an alpha is $\frac{1}{4}$ that of a proton.
 - Use proton at 1.92 MeV
- Approximate tissue as water and find proton range from a table.
 - 2 MeV, $R_p = 0.007 \text{ g cm}^{-2}$
- Units are density adjusted.
 - $R_\alpha = 0.007 \text{ cm}$
- Alpha can't penetrate the skin.

Electron Interactions

- Electrons share the same interactions as protons.
 - Coulomb interactions with atomic electrons
 - Low mass changes result
- Electrons also have stopping radiation: bremsstrahlung
- Positrons at low energy can annihilate.



Beta Collisions

- There are a key differences between betas and heavy ions in matter.
 - A large fractional energy change
 - Indistinguishable β^- from e in quantum collision
- Bethe formula is modified for betas.

$$\left(-\frac{dE}{dx}\right)_{\text{col}}^{\pm} = \frac{4\pi n k^2 e^4}{mc^2 \beta^2} \left[\ln \frac{2mc^2 \tau \sqrt{\tau+2}}{\sqrt{2}I} - F^{\pm}(\beta) \right] \quad \tau = \frac{K}{mc^2}$$

$$F^{-}(\beta) = \frac{1-\beta^2}{2} \left[1 + \frac{\tau^2}{8} - (2\tau+1) \ln 2 \right]$$

$$F^{+}(\beta) = \ln 2 - \frac{\beta^2}{24} \left[23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right]$$

Radiative Stopping

- The energy loss due to bremsstrahlung is based classical electromagnetism.
 - High energy
 - High absorber mass

$$\left(-\frac{dE}{dx}\right)_{\text{rad}} = \frac{4nk^2Z(Z+1)e^4\tau}{137mc^2} \left[\ln 2\tau - \frac{1}{3} \right]$$

- There is an approximate relation.

$$\left(\frac{dE}{dx}\right)_{\text{rad}} = \frac{KZ}{700} \left(\frac{dE}{dx}\right)_{\text{col}}^{\pm}$$

-(dE/dx)/ρ in water

Energy	MeV cm ² g ⁻¹	
	col	rad
1 keV	126	0
10 keV	23.2	0
100 keV	4.2	0
1 MeV	1.87	0.017
10 MeV	2.00	0.183
100 MeV	2.20	2.40
1 GeV	2.40	26.3

Beta Range

- The range of betas in matter depends on the total dE/dx .
 - Energy dependent
 - Material dependent
- Like other measures, it is often scaled to the density.

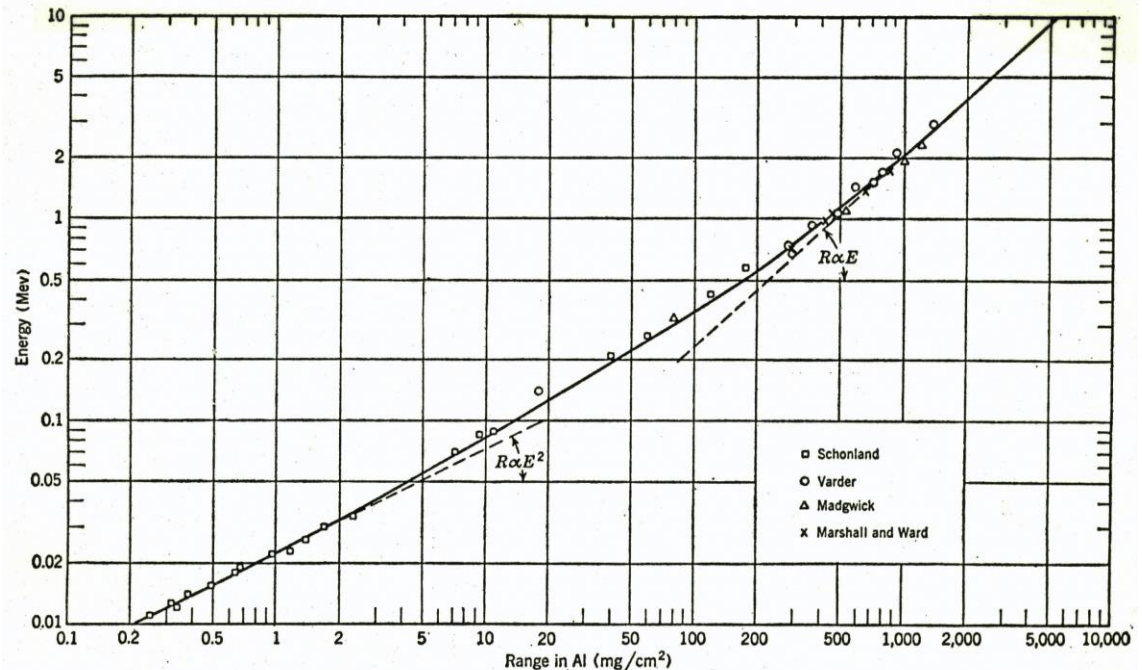


Fig. 3.3 Empirical range-energy relationship for electrons absorbed in aluminum. Experimental values by several observers (S16, V4, M4, M17) on monoenergetic electrons are shown. For monoenergetic electrons, the range coordinate refers to the extrapolated range R_0 of Fig. 3.2. For continuous β -ray spectra the energy coordinate refers to the end-point energy E_{max} and the range coordinate becomes the maximum range R_m of Fig. 3.4. The smooth curve represents the empirical relationship, Eqs. (3.3) and (3.4), developed by Katz and Penfold (K7).

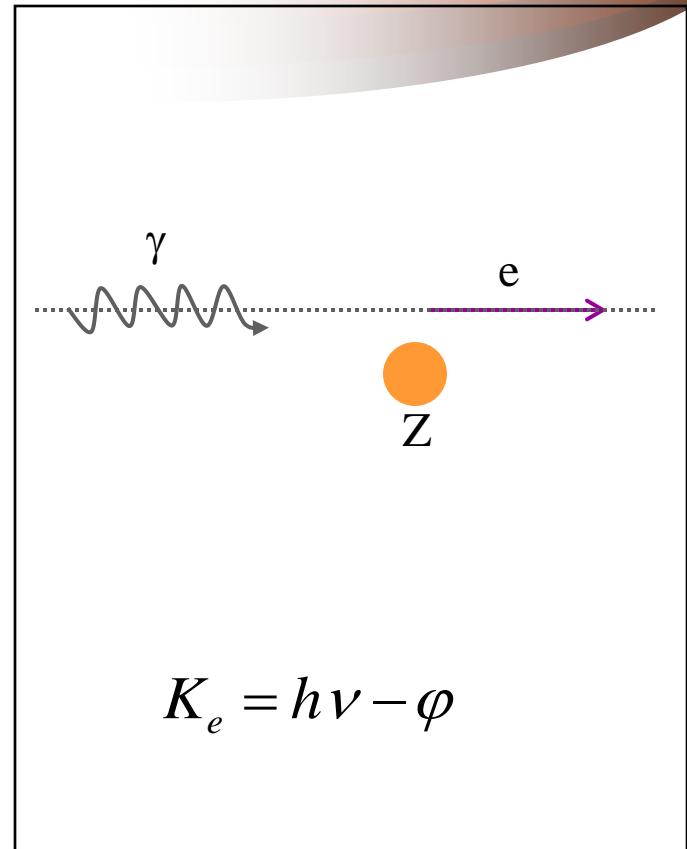
Photon Interactions



- High energy photons interact with electrons.
 - Photoelectric effect
 - Compton effect
- They also indirectly interact with nuclei.
 - Pair production

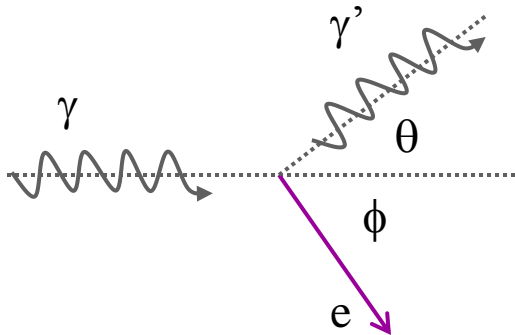
Photoelectric Effect

- A photon can eject an electron from an atom.
 - Photon is absorbed
 - Minimum energy needed for interaction.
 - Cross section decreases at high energy



Compton Effect

- Photons scattering from atomic electrons are described by the Compton effect.
 - Conservation of energy and momentum



$$h\nu + mc^2 = h\nu' + E'$$

$$\frac{h\nu}{c} = \frac{h\nu'}{c} \cos \theta + P' \cos \phi$$

$$\frac{h\nu'}{c} \sin \theta = P' \sin \phi$$

$$h\nu' = \frac{h\nu}{1 + \frac{h\nu}{mc^2} (1 - \cos \theta)}$$

Compton Energy

- The frequency shift is independent of energy.
- The energy of the photon depends on the angle.
 - Max at 180°
- Recoil angle for electron related to photon energy transfer
 - Small $\theta \rightarrow$ cot large
 - Recoil near 90°

$$\Delta\lambda = \frac{h}{mc} (1 - \cos \theta)$$

$$K = \frac{h\nu(1 - \cos \theta)}{mc^2 / h\nu + 1 - \cos \theta}$$

$$\cot \frac{\theta}{2} = \left(1 + \frac{h\nu}{mc^2} \right) \tan \phi$$

Compton Cross Section

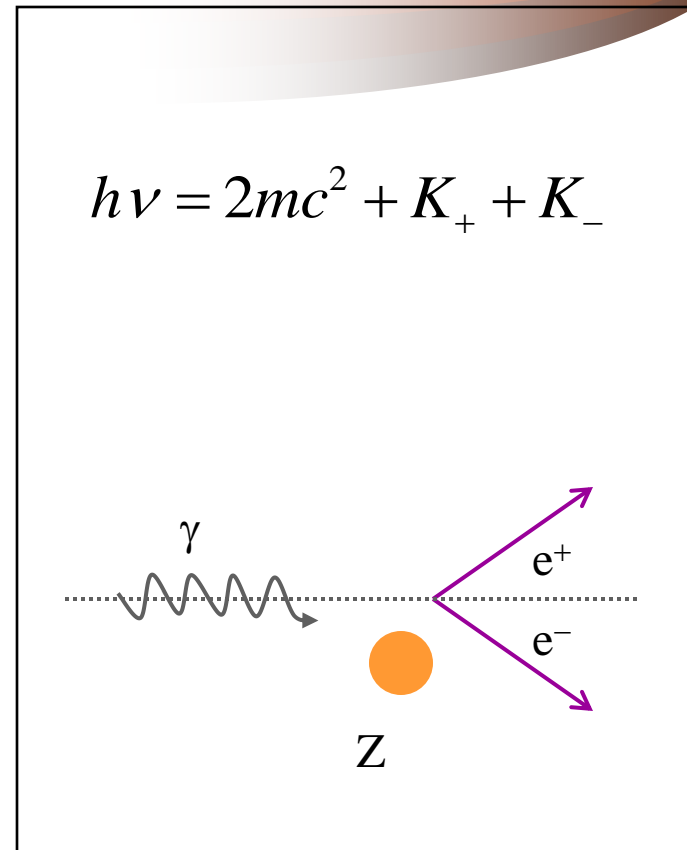
- Differential cross section can be derived quantum mechanically.
 - Klein-Nishina
 - Scattering of photon on one electron
 - Units $\text{m}^2 \text{sr}^{-1}$
- Integrate to get cross section per electron
 - Multiply by electron density
 - Units m^{-1}

$$\frac{d\sigma}{d\Omega} = \frac{k^2 e^4}{2m^2 c^4} \left(\frac{\nu'}{\nu}\right)^2 \left(\frac{\nu}{\nu'} + \frac{\nu'}{\nu} - \sin^2 \theta\right)$$

$$\sigma = 2\pi n \int \frac{d\sigma}{d\Omega} \sin \theta d\theta$$

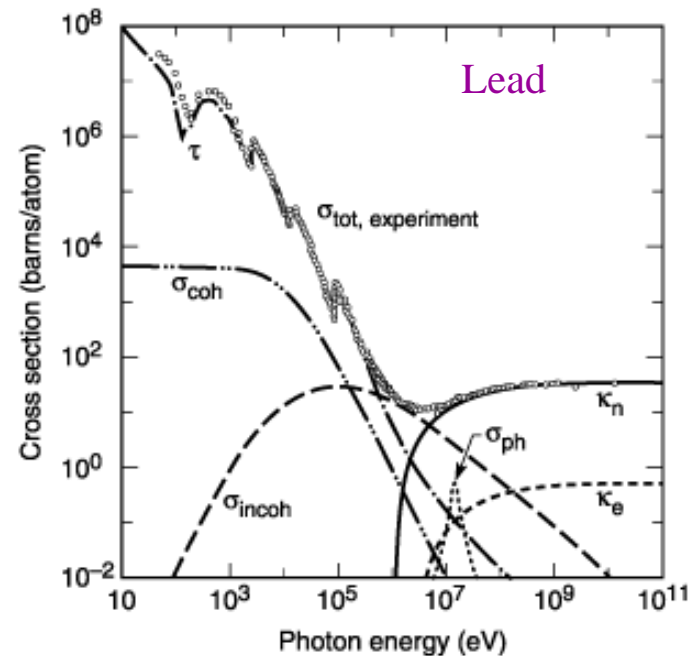
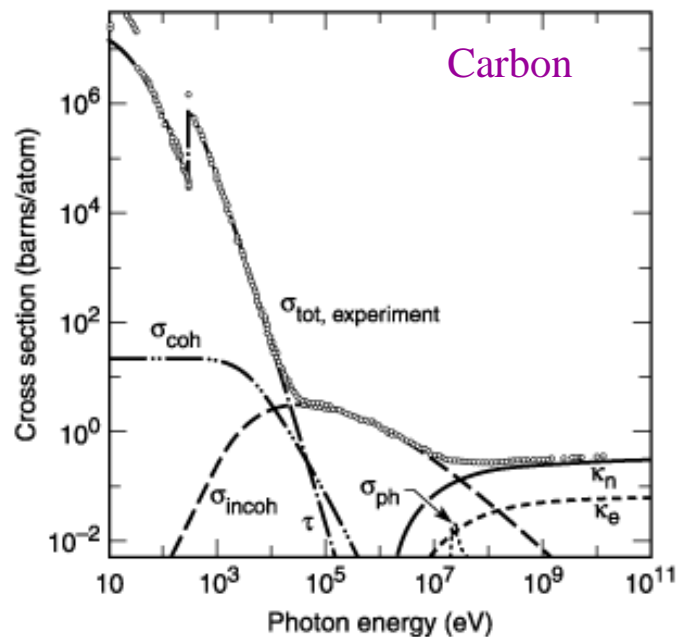
Pair Production

- Photons above twice the electron rest mass energy can create an electron-positron pair.
 - Minimum $\lambda = 0.012 \text{ \AA}$
- The nucleus is involved for momentum conservation.
 - Probability increases with Z
- This is a dominant effect at high energy.



Total Photon Cross Section

- Photon cross sections are the sum of all effects.
 - Photoelectric τ , Compton σ_{incoh} , pair κ



J. H. Hubbell (1980)